

Study of an Expanding Bubble in the Rayleigh Model

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Abstract: In this study we investigate the expanding bubble in the rayleigh model. A bubble expands adiabatically in an incompressible, inviscid liquid. The variation of its radius R with time is given by the Rayleigh's equation. We find that the bubble is stable at the equilibrium point in this model.

Keywords: Bubble, Rayleigh's model and stability

I. INTRODUCTION :

When a bubble expands adiabatically in an inviscid liquid at rest at infinity, the variation of its radius R with time is given by Rayleigh's equation, which is a highly non-linear equation. This equation has been generalized including viscosity of surrounding liquid and surface tension by Noltingk and Neppiras and Poritsky. Progress in the study of this equation is generally expected numerically. However, when the liquid outside the bubble is inviscid, Rayleigh's modified equation can be transformed suitably and we can prove, analytically, that the expanding bubble in this model is stable at the equilibrium point . .

II. MATHEMATICAL ANALYSIS :

As the bubble expands adiabatically, its radius R is given by Rayleigh's equation, modified by surface tension σ . This equation can be written as –

$$R \frac{d^2 R}{dt^2} + \frac{3}{2} \left(\frac{dR}{dt} \right)^2 + \frac{1}{\rho} \left\{ p_e - p_{\infty} \left(\frac{R_0}{R} \right)^{\gamma} + \frac{2\sigma}{R} \right\} = 0, \tag{1}$$

Where ρ is the density of the outside inviscid liquid, ρ_{∞} and R_0 are the gas pressure and radius of the bubble initially, p_e is the pressure, taken as constant, in the liquid at a large distance from the bubble and γ is the ratio of two specific heats of the gas. Then find that equation (1) can be written in the form-

$$\frac{d^2}{dt^2} (R^{5/2}) + \frac{5R^{1/2}}{2\rho} \left\{ p_e - p_{\infty} \left(\frac{R_0}{R} \right)^{\gamma} + \frac{2\sigma}{R} \right\} = 0. \tag{2}$$

Defining

$$r = R^{5/2}$$

and taking $\gamma = 4/3$ for simplicity, we finally get from

(2) the equation

$$\frac{d^2 r}{dt^2} + \frac{5}{2\rho} r^{1/5} \left\{ p_e - p_{\infty} \left(\frac{r_0}{r} \right)^{8/5} + 2\sigma r^{-2/5} \right\} = 0, \tag{3}$$

where $r_0 = R_0^{5/2}$

We use ρ_{∞} and U_0 as characteristic pressure and speed respectively and r_0 as the characteristic value of r and T_0 as that of time, where $T_0 = R_0/U_0$. Then define the dimensionless quantities r' , t' and $p'e$ as

$$R' = r/r_0, t' = t/T_0, p'e = p_e/p_{\infty}$$

We may now write eq. (3) as

$$\frac{d^2 r'}{dt'^2} = \frac{-5p_{\infty}}{2\rho U_0^2} \left\{ p'e r'^{1/5} + \frac{2\sigma}{p_{\infty} R_0} r'^{-1/5} - r'^{-7/5} \right\}. \tag{4}$$

Omitting the dashes from r' , t' and $p'e$ and using from now on these undashed symbols for the corresponding quantities, then find that eq (4) can be written in dimensionless form as

$$\frac{d^2 r}{dt^2} = \frac{-5p_{\infty}}{2\rho U_0^2} \left\{ p_e r^{1/5} + \frac{2\sigma}{p_{\infty} R_0} r^{-1/5} - r^{-7/5} \right\}. \tag{5}$$

III. CONCLUSION

We find that F (x) vanishes at the equilibrium point. Also H and its partial derivatives are continuous at all points except when $x = 0$. If α is a positive real root, then shows that $x = \alpha$ and $y = 0$ is the equilibrium point. Let us choose C in such a way that at this equilibrium point, the value of H vanishes. Which is always satisfied as x is real and positive. Therefore, H is minimum at the equilibrium point $x = \alpha$, $y = 0$, but H vanishes at $(\alpha, 0)$. Thus H is always positive near $(\alpha, 0)$ and is therefore positive definite in the neighbourhood of the equilibrium point.

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